

# The exact Hohenberg-Kohn functional for a lattice model

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MAX-PLANCK-GESELLSCHAFT

## Introduction

For a discretized soft-Coulomb lattice model, we investigate the exact solution of the many-body Schrödinger equation in Fock space. Using quadratic optimization with quadratic constraints, or alternatively exact diagonalization, we explicitly construct the exact Hohenberg-Kohn functional and the mapping from densities to wavefunctions. We analyze the resulting exact Hohenberg-Kohn functional and draw conclusions for the construction of approximate functionals.

## Levy-Lieb constraint search (M. Levy 1979 [1], E. Lieb 1983 [2])

Expand eigenfunctions in a complete basis set (energy eigenfunctions, Slater-Determinants, etc.)

$$|\Psi[n]\rangle = \sum_{j=1}^M \alpha_j[n] |\phi_j\rangle, \text{ M number of sites}$$

Hohenberg-Kohn functional

$$F_{\text{HK}}(\alpha_1, \dots, \alpha_M)[n] = \min_{\Psi \rightarrow n} \langle \Psi[n] | \hat{T} + \hat{W} | \Psi[n] \rangle$$

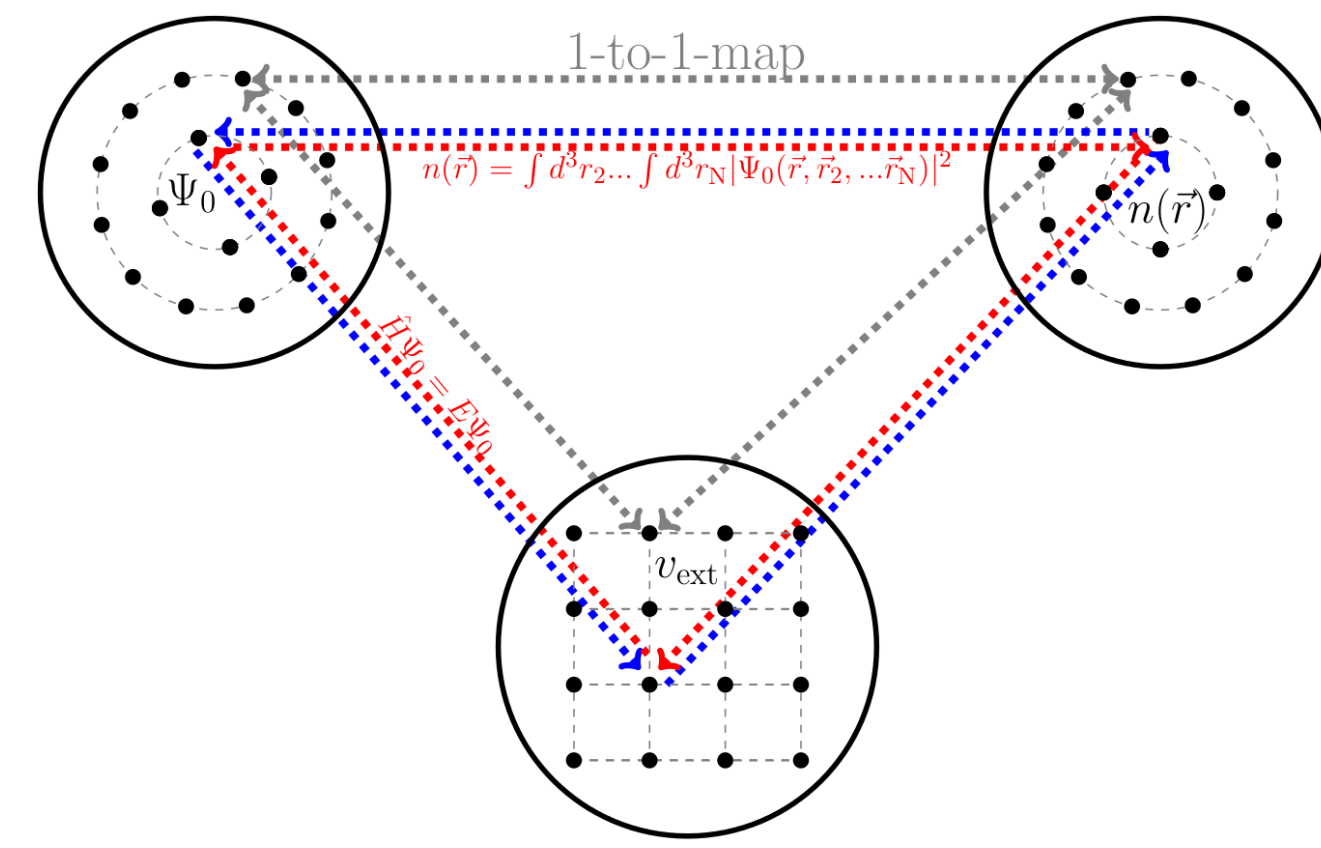
$$= \min_{\Psi \rightarrow n} \sum_{j,k=1}^M \alpha_j^* \alpha_k [n] \langle \phi_j | \hat{T} + \hat{W} | \phi_k \rangle$$

## Two-site soft-Coulomb model

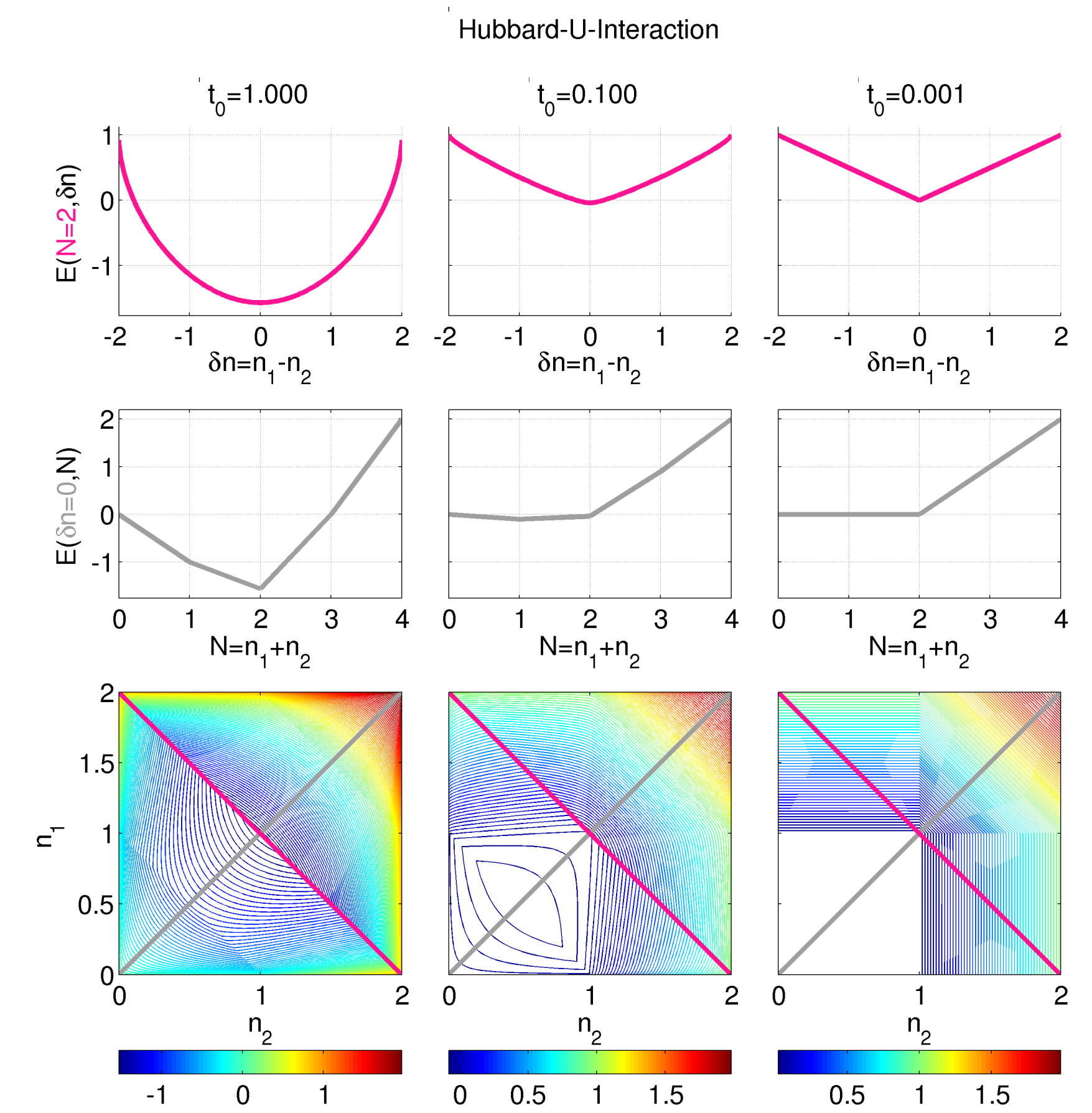
$$\hat{H} = \hat{T} + \hat{W} + \hat{V}, \quad \hat{T} = -t_0 \sum_{l,\sigma} (\hat{c}_{l,\sigma}^\dagger \hat{c}_{l+1,\sigma} + \hat{c}_{l+1,\sigma}^\dagger \hat{c}_{l,\sigma}) + 2t_0 \sum_{l,\sigma} \hat{n}_{l,\sigma}, \quad t_0 = \frac{\hbar^2}{2m_0 \Delta^2},$$

$$\hat{W}_{\text{H}} = U \sum_{l,\sigma} \hat{n}_{l,\sigma} \hat{n}_{l,\sigma}, \quad \hat{W}_{\text{SC}} = \sum_{l,m,\sigma,\sigma'} \frac{e^2 \hat{c}_{l,\sigma}^\dagger \hat{c}_{m,\sigma'}^\dagger \hat{c}_{m,\sigma'} \hat{c}_{l,\sigma}}{2\sqrt{(l-m\Delta)^2+1}}, \quad \hat{V} = \sum_{l,\sigma} \hat{n}_{l,\sigma} \cdot v_{l,\sigma}, \quad \hat{n} = \sum_{j,\sigma} \hat{c}_{j,\sigma}^\dagger \hat{c}_{j,\sigma}$$

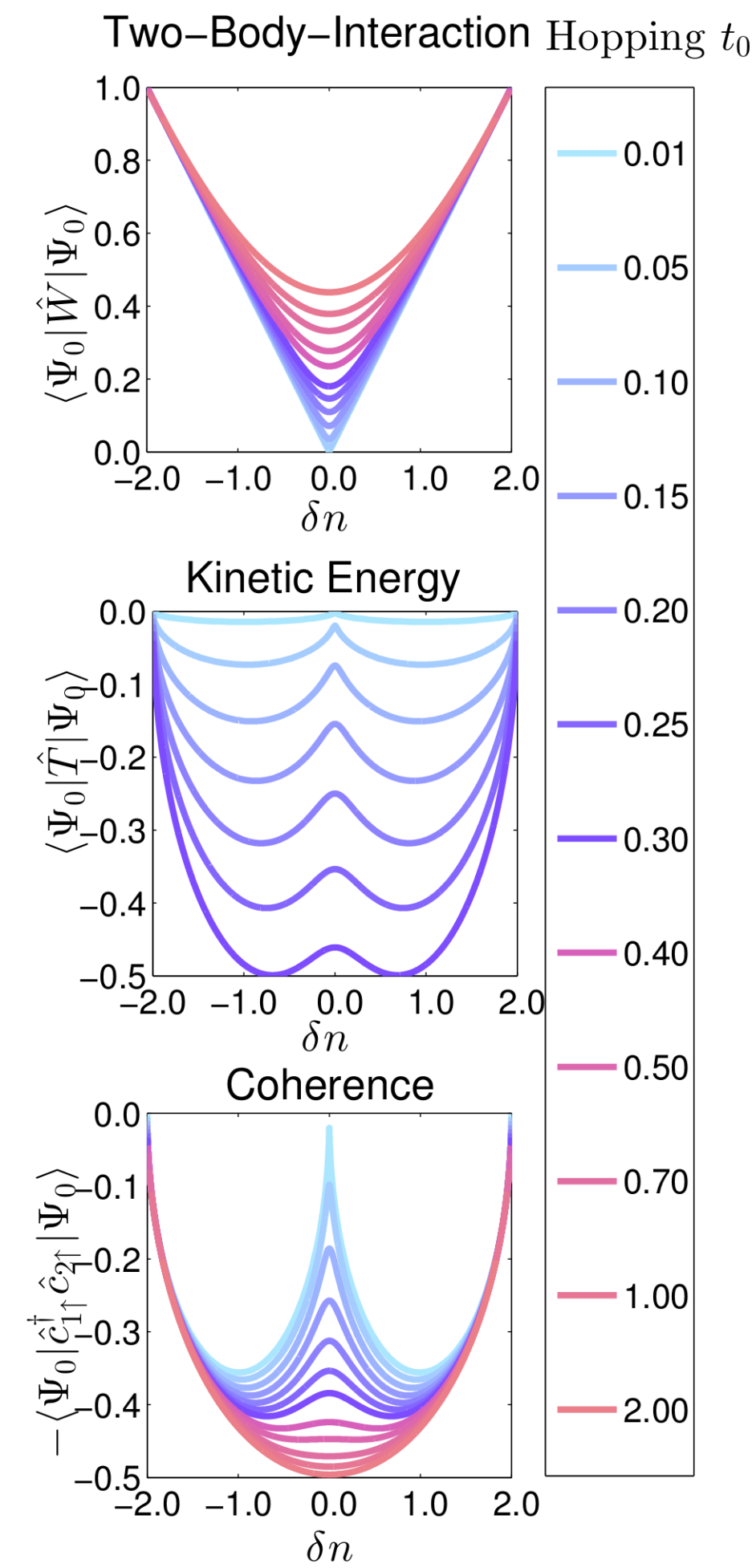
We consider different particle numbers by including a chemical potential  $\mu$ . Spacing  $\Delta$



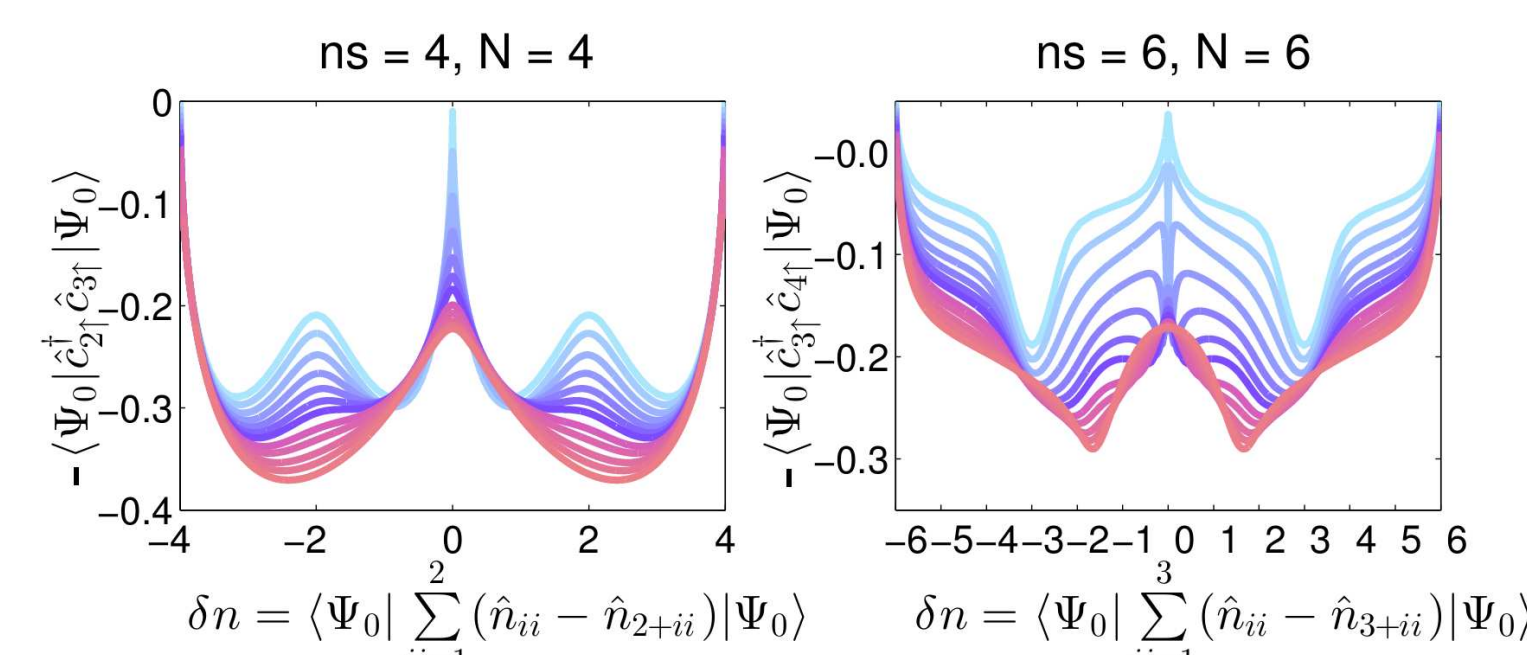
## Exact Hohenberg-Kohn-Functional



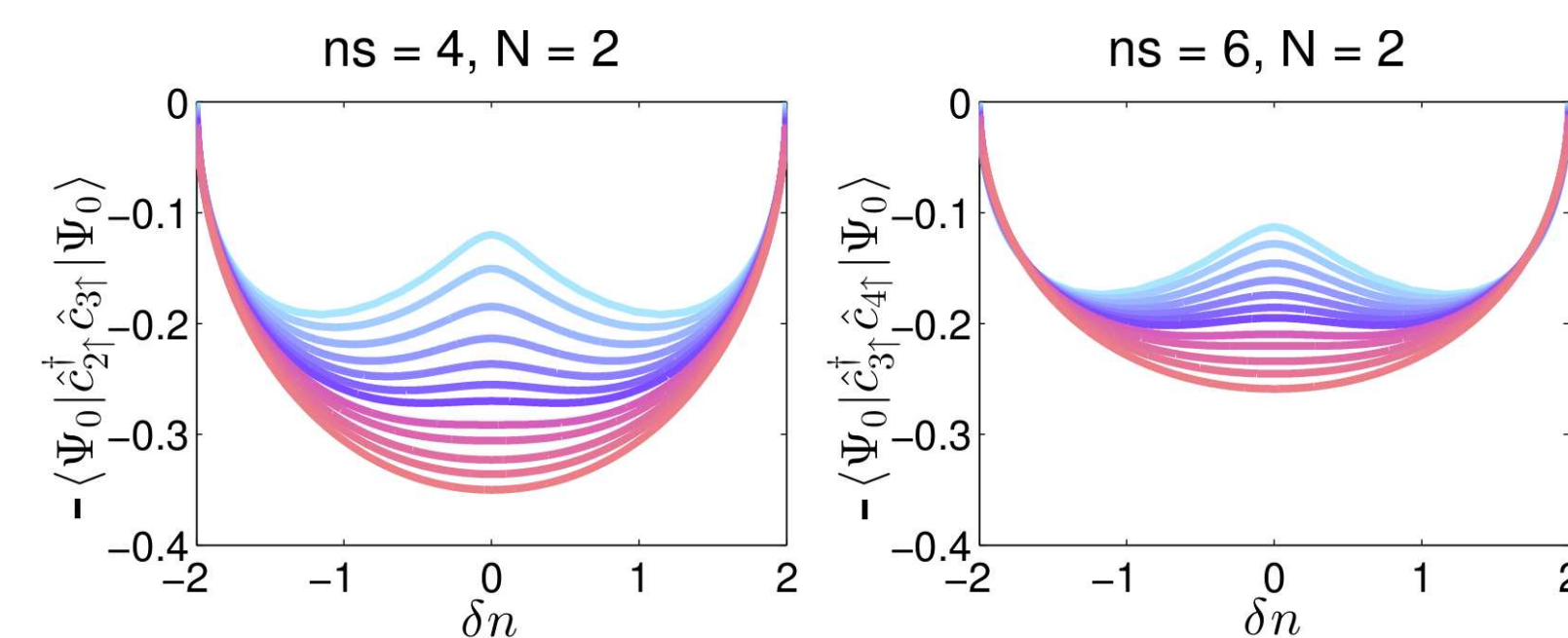
## Exact Coherences



## Coherences for ns = 4 and ns = 6 (half-filling)



## Coherences for ns = 4 and ns = 6 (N=2)



## Soft-Coulomb molecules in 1D

Hamiltonian

$$\hat{H}(\alpha) = \hat{T} + \hat{W} + \hat{V}(\alpha)$$

$$\hat{T} = \sum_{j=1}^2 -\frac{\hbar^2}{2m} \frac{d^2}{dx_j^2}, \quad \hat{W} = \frac{1}{2} \sum_{i \neq j} \frac{1}{\sqrt{(x_i - x_j)^2 + 1}}$$

$$\hat{V}(\alpha) = \sum_{j=1}^2 \frac{Z_1(\alpha)}{\sqrt{(x_j - d)^2 + 1}} + \frac{Z_2(\alpha)}{\sqrt{(x_j + d)^2 + 1}}$$

$$Z_1(\alpha) = -\alpha, \quad Z_2(\alpha) = -(2 - \alpha), \quad \alpha \in [0, 2], \quad d = 3, 8 \text{ Bohr}$$

Exact Kohn-Sham potential for two electrons in spin singlet configuration (Helbig et al. 2009 [3])

$$v_{\text{KS}}(x) = \frac{1}{2} \frac{\nabla^2 \sqrt{n(x)}}{\sqrt{n(x)}} + \epsilon_1$$

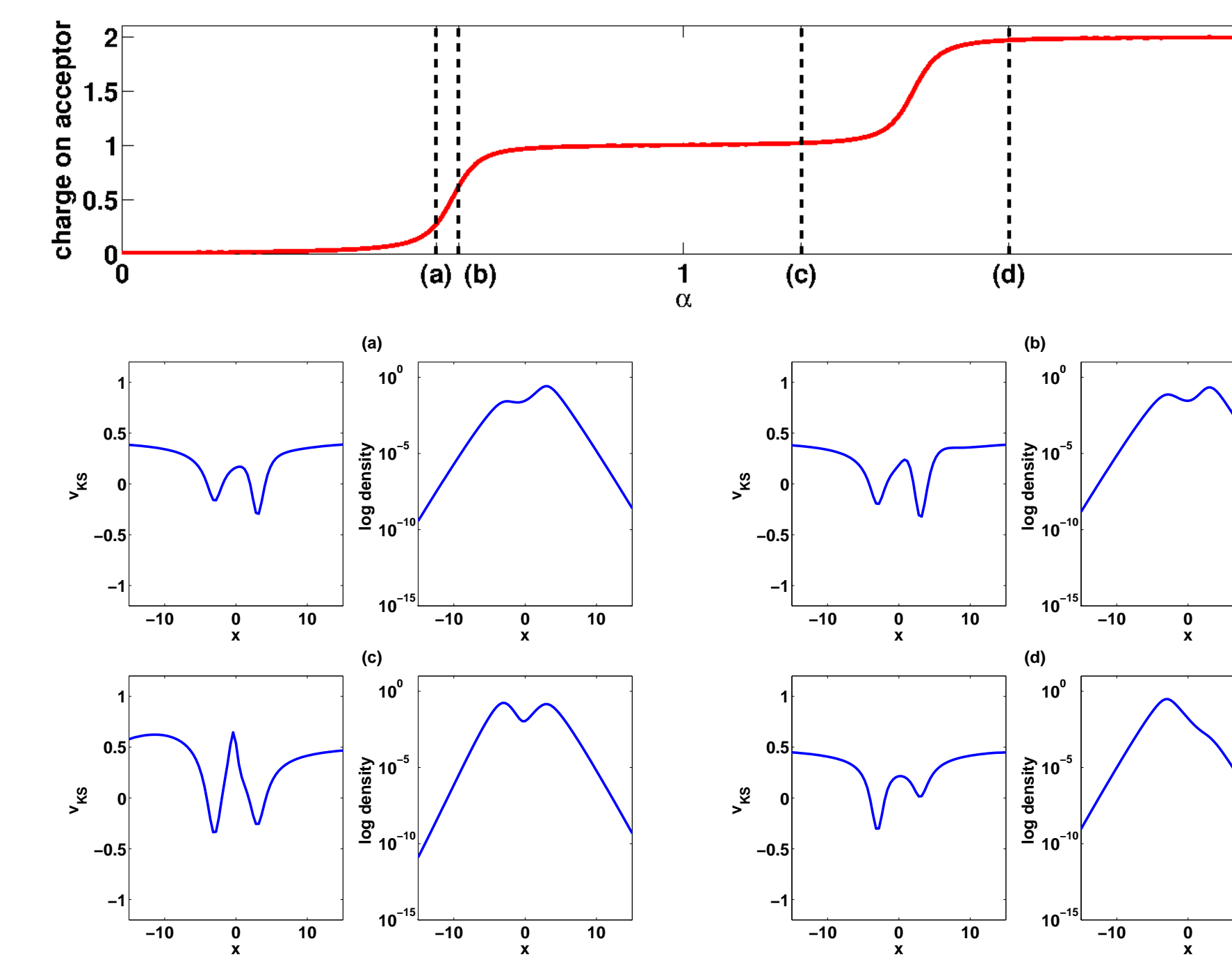
Exact solution of static two-electron Schrödinger equation with octopus (A. Castro et al. [4])

$$\hat{H}(\alpha) \Psi_j(\alpha) = E_j(\alpha) \Psi_j(\alpha)$$

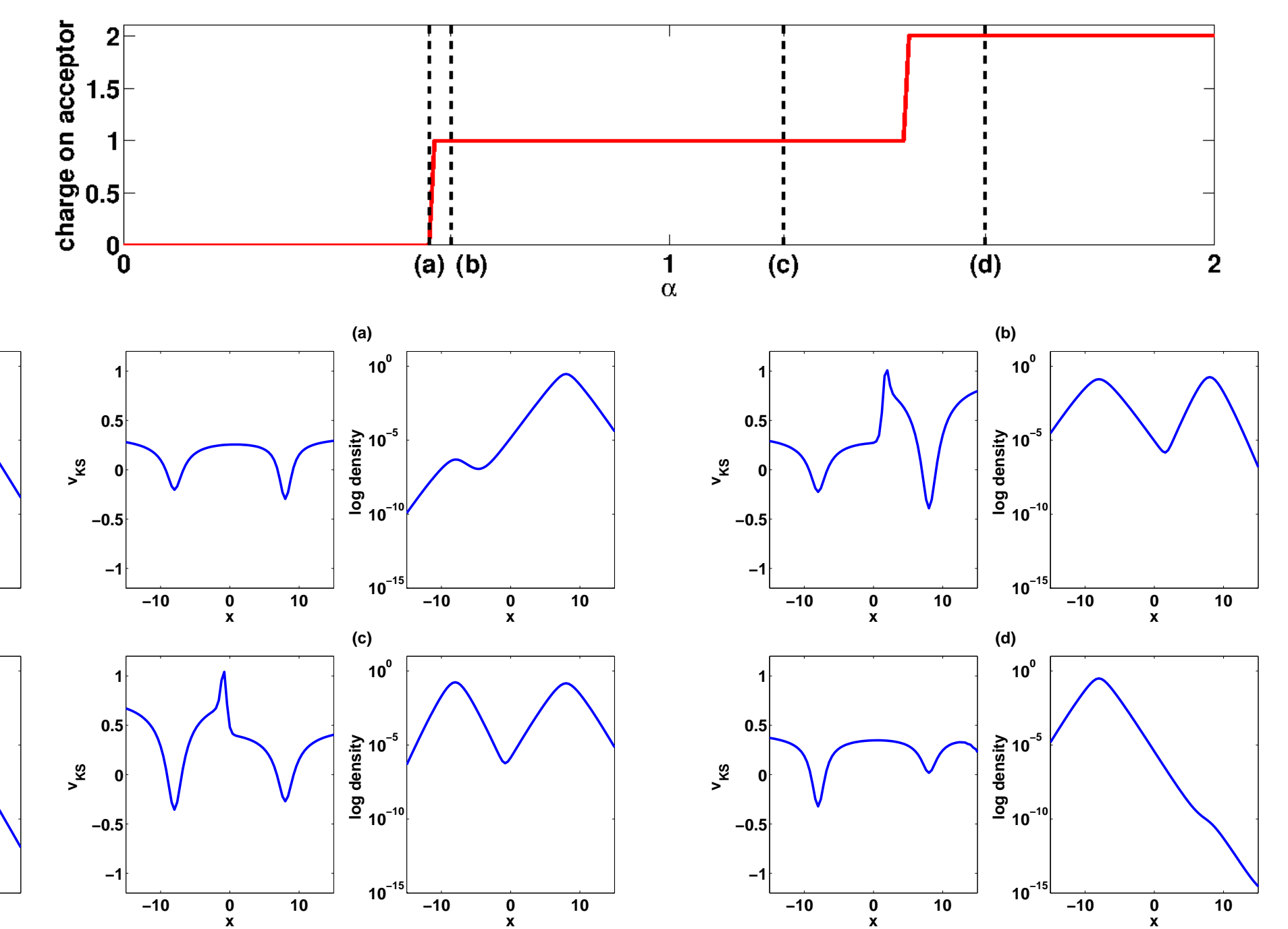
$$n(x) = \langle \Psi | \hat{n}(x) | \Psi \rangle$$

$$\hat{n}(x) = \sum_j \delta(x - x_j)$$

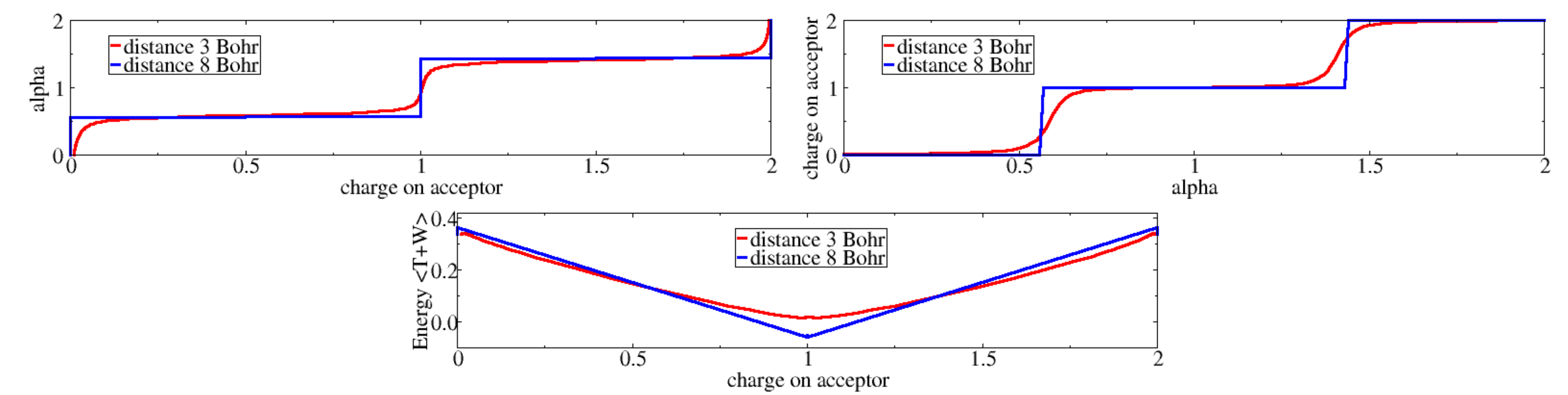
## High-density limit (small distance d = 3 Bohr)



## Low-density limit (large distance d = 8 Bohr)



## Softened intra-system derivative discontinuity



## Conclusion & Outlook

- The exact Hohenberg-Kohn functional shows a softened intra-system derivative discontinuity in the low-density limit.
- Expectation values of operators are affected by the softened intra-system derivative discontinuity.
- We observe softened intra-system derivative discontinuity also for soft-Coulomb molecules in 1D.
- We currently develop an approximate functional which incorporates the intra-system derivative discontinuity.

## References

- [1] M. Levy, Proc. Natl Acad. Sci. USA 76 6062 (1979)
- [2] E. Lieb, Int. J. Quantum Chem. 24 24377 (1983)
- [3] N. Helbig, I.V. Tokatly, A. Rubio, Journal of Chemical Physics 131, 224105 (2009)
- [4] A. Castro, H. Appel, Micael Oliveira, C.A. Rozzi, X. Andrade, F. Lorenzen, M.A.L. Marques, E.K.U. Gross, and A. Rubio, Phys. Stat. Sol. B 243 2465-2488 (2006)
- [5] John P. Perdew, Robert G. Parr, Mel Levy, Jose L. Balduz, Jr., Phys. Rev. Lett. 49, 16911694 (1982)